

## On almost hyperHermitian structures on Riemannian manifolds and tangent bundles

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**Abstract:** Some results concerning almost hyperHermitian structures are considered, using the notions of the canonical connection and the second fundamental tensor field  $h$  of a structure on a Riemannian manifold which were introduced by the second author.

With the help of any metric connection  $\tilde{\nabla}$  on an almost Hermitian manifold  $M$  an almost hyperHermitian structure can be constructed in the defined way on the tangent bundle  $TM$ . A similar construction was considered in [6], [7]. This structure includes two basic anticommutative almost Hermitian structures for which the second fundamental tensor fields  $h^1$  and  $h^2$  are computed. It allows us to consider various classes of almost hyperHermitian structures on  $TM$ . In particular, there exists an infinite-dimensional set of almost hyperHermitian structures on  $TTM$  where  $M$  is any Riemannian manifold.

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### 1 Some remarks on almost hyperHermitian structures

$1^0$ . Let  $(M, J, g)$  be an almost Hermitian manifold i.e.  $J^2 = -I$  and  $g(JX, JY) = g(X, Y)$  for  $X, Y \in \chi(M)$ , where  $g$  is a fixed Riemannian metric on  $M$ . For any Riemannian metric  $\tilde{g}$  on  $M$  such a metric  $g$  is defined by the formula

$$g(X, Y) = \frac{1}{2}(\tilde{g}(X, Y) + \tilde{g}(JX, JY)), \quad X, Y \in \chi(M).$$

Let  $\nabla$  be the Riemannian connection of the metric  $g$ . Then one can define a connection  $\bar{\nabla}$  on  $M$  by

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$$\bar{\nabla}_X Y = \frac{1}{2}(\nabla_X Y - J\nabla_X JY) = \nabla_X Y + \frac{1}{2}\nabla_X(J)JY, X, Y \in \chi(M). \tag{1}$$

The connection  $\bar{\nabla}$  is called the canonical connection of the pair  $(J, g)$ , or more precisely of the corresponding  $G$ -structure, where  $G = U(n)$ , [2]. In particular,  $\bar{\nabla}g = 0$ ,  $\bar{\nabla}J = 0$ .

The tensor field  $h$  is called the second fundamental tensor field of the pair  $(J, g)$  [2], where

$$h_X Y = \nabla_X Y - \bar{\nabla}_X Y = -\frac{1}{2}\nabla_X(J)JY = \frac{1}{2}(\nabla_X Y + J\nabla_X JY), X, Y \in \chi(M); \tag{2}$$

$$h_{XYZ} = g(h_X Y, Z) = -h_{XZY}. \tag{3}$$

In particular, the classification given in [3] can be rewritten in terms of the tensor field  $h$ , [2]. Let  $\dim M \geq 6$  and  $2\beta(X) = \delta\Phi(JX)$ , where  $\Phi(X, Y) = g(JX, Y)$ . Then we have

Class	Defining condition
<b>K</b>	$h = 0$
<b>U<sub>1</sub>=NK</b>	$h_X X = 0$
<b>U<sub>2</sub>=AK</b>	$\sigma h_{XYZ} = 0$
<b>U<sub>3</sub>=SK∩H</b>	$h_{XYZ} - h_{JXJYZ} = \beta(Z) = 0$
<b>U<sub>4</sub></b>	$h_{XYZ} = \frac{1}{2(n-1)}[\langle X, Y \rangle \beta(Z) - \langle X, Z \rangle \beta(Y) - \langle X, JY \rangle \beta(JZ) + \langle X, JZ \rangle \beta(JY)]$
<b>U<sub>1</sub>⊕U<sub>2</sub>=QK</b>	$h_{XYJZ} = h_{JXYZ}$
<b>U<sub>3</sub>⊕U<sub>4</sub>=H</b>	$N(J) = 0$ or $h_{XYJZ} = -h_{JXYZ}$
<b>U<sub>1</sub>⊕U<sub>3</sub></b>	$h_{XXY} - h_{JXJXY} = \beta(Z) = 0$
<b>U<sub>2</sub>⊕U<sub>4</sub></b>	$\sigma[h_{XYJZ} - \frac{1}{(n-1)}\langle JX, Y \rangle \beta(Z)] = 0$
<b>U<sub>1</sub>⊕U<sub>4</sub></b>	$h_{XXY} = -\frac{1}{2(n-1)}[\langle X, Y \rangle \beta(X) - \ X\ ^2 \beta(Y) - \langle X, JY \rangle \beta(JX)]$
<b>U<sub>2</sub>⊕U<sub>3</sub></b>	$\sigma[h_{XYJZ} + h_{JXYZ}] = \beta(Z) = 0$
<b>U<sub>1</sub>⊕U<sub>2</sub>⊕U<sub>3</sub>=SK</b>	$\beta = 0$
<b>U<sub>1</sub>⊕U<sub>2</sub>⊕U<sub>4</sub></b>	$h_{XYJZ} - h_{JXYZ} = \frac{1}{(n-1)}[\langle X, Y \rangle \beta(JZ) - \langle X, Z \rangle \beta(JY) + \langle X, JY \rangle \beta(Z) - \langle X, JZ \rangle \beta(Y)]$
<b>U<sub>1</sub>⊕U<sub>3</sub>⊕U<sub>4</sub></b>	$h_{XJXY} + h_{JXXY} = 0$
<b>U<sub>2</sub>⊕U<sub>3</sub>⊕U<sub>4</sub></b>	$\sigma[h_{XYJZ} + h_{JXYZ}] = 0$
<b>U</b>	No condition

Table 1

2<sup>0</sup>. We consider an almost hyperHermitian structure (ahHs) on a manifold  $M$  consisting of  $(J_1, J_2, J_3, g)$ , where  $J_i^2 = -I, J_1J_2 = -J_2J_1 = J_3, g(J_iX, J_iY) = g(X, Y), i = 1, 2, 3$ .

For any Riemannian metric  $\tilde{g}$  such a metric  $g$  can be defined by the formula  $g(X, Y) = \frac{1}{4}(\tilde{g}(X, Y) + \tilde{g}(J_1X, J_1Y) + \tilde{g}(J_2X, J_2Y) + \tilde{g}(J_3X, J_3Y)), X, Y \in \chi(M)$ .

If  $\nabla$  is the Riemannian connection of the metric  $g$ , then the canonical connection  $\bar{\nabla}$  in the sense of [2] of the ahHs has the following form

$$\bar{\nabla}_X Y = \frac{1}{4}(\nabla_X Y - J_1 \nabla_X J_1 Y - J_2 \nabla_X J_2 Y - J_3 \nabla_X J_3 Y), X, Y \in \chi(M). \tag{4}$$

In particular,  $\bar{\nabla}g = 0, \bar{\nabla}J_i = 0, i = 1, 2, 3$ .

**Proposition 1.1.** *Let  $(M, J_1, g)$  be a Kaehlerian manifold i.e.  $\nabla J_1 = 0$  on  $M$ . Then the connection given by (4) coincides with those defined by (1) for  $(M, J_2, g)$  and  $(M, J_3, g)$ . In particular, the second fundamental tensor fields of  $(M, J_2, g)$  and  $(M, J_3, g)$  are the same.*

**Proof.**

$$\begin{aligned} \bar{\nabla}_X Y &= \frac{1}{4}(\nabla_X Y - J_1^2 \nabla_X Y - J_2 \nabla_X J_2 Y - J_1 J_2 \nabla_X J_1 J_2 Y) = \\ &= \frac{1}{4}(2\nabla_X Y - J_2 \nabla_X J_2 Y - J_1 J_2 J_1 \nabla_X J_2 Y) = \\ &= \frac{1}{2}(\nabla_X Y - J_2 \nabla_X J_2 Y) = \\ &= \frac{1}{2}(\nabla_X Y - J_3 J_1 \nabla_X J_3 J_1 Y) = \\ &= \frac{1}{2}(\nabla_X Y - J_3 \nabla_X J_3 Y). \end{aligned}$$

□

To illustrate the situation described in *proposition 1.1* we consider the following example.

**Example 1.2.** Let  $(M, J, g)$  be an almost Hermitian manifold and let  $(J, g)$  belong to one of the classes in *Table 1*,  $\dim M = 4n$ . Further, define orthonormal vector fields  $X_1, \dots, X_{2n}, JX_1, \dots, JX_{2n}$  on some open neighborhood  $U$  of a point  $p \in M$ . Assuming  $J_1 = J$  on  $U$ , we get an almost Hermitian manifold  $(U, J_1, g)$  which also belongs to the corresponding class. Then, we can define  $J_2$  by the following equalities

$$\begin{aligned} J_2 X_1 &= X_{n+1}, J_2 X_2 = X_{n+2}, \dots, J_2 X_n = X_{2n}; \\ J_2 X_{n+1} &= -X_1, J_2 X_{n+2} = -X_2, \dots, J_2 X_{2n} = -X_n; \\ J_2(J_1 X_1) &= -J_1 X_{n+1}, J_2(J_1 X_2) = -J_1 X_{n+2}, \dots, J_2(J_1 X_n) = -J_1 X_{2n}; \\ J_2(J_1 X_{n+1}) &= J_1 X_1, J_2(J_1 X_{n+2}) = J_1 X_2, \dots, J_2(J_1 X_{2n}) = J_1 X_n. \end{aligned}$$

It is clear that  $J_1 J_2 = -J_2 J_1, J_1^2 = J_2^2 = -I$  and  $g(J_i X, J_i Y) = g(X, Y), i = 1, 2, 3$  on  $U$ , where  $J_3 = J_1 J_2, X, Y \in \chi(U)$ .

If  $(M, J, g)$  is a Kaehlerian manifold (class **K**) and  $\dim M = 4n$  then we obtain the situation given in *proposition 1.1* on the almost hyperHermitian manifold

$(U, J_1, J_2, J_3, g)$ .

One can find examples of the structures from Table 1 in [2], [3].

To get an almost Hermitian manifold of dimension  $4n$  we can take a Riemannian product  $M \times M$ .

**Problem 1.3.** Let  $(U_\alpha, U_\beta)$  be any pair of the classes from Table 1 ( $\alpha, \beta = 0, \dots, 15$ ). Can one construct such examples of ahHs  $(J_1, J_2, J_3, g)$  on manifolds such that  $(J_1, g)$  belongs to the class  $U_\alpha$  and  $(J_2, g)$  belongs to the class  $U_\beta$ ? The cases  $(K, K)$  and  $(U, U)$  are easily illustrated.

**3<sup>0</sup>. Definition 1.4.** [5] *A connected Riemannian manifold  $(M, g)$  with a family of local isometries  $\{s_x: x \in M\}$  is called a locally  $k$ -symmetric Riemannian space ( $k - s. l. R. s.$ ) if the following axioms are fulfilled:*

- (a)  $s_x(x) = x$  and  $x$  is the isolated fixed point of the local symmetry  $s_x$ ;
- (b) the tensor field  $S: S_x = (s_x)_{*x}$  is smooth and invariant under any local isometry  $s_x$ ;
- (c)  $S^k = I$  and  $k$  is the least of such positive integers.

If  $M$  is a  $k$ -s. l. R. s. then the unique canonical connection  $\tilde{\nabla}$  can be defined by the following formula (see [2])

$$\tilde{\nabla}_X Y = \nabla_X Y - \frac{1}{k} \sum_{j=1}^{k-1} \nabla_X (S^j) S^{k-j} Y = \frac{1}{k} \sum_{j=0}^{k-1} S^j \nabla_X S^{k-j} Y, X, Y \in \chi(M). \tag{5}$$

Further, we have  $\tilde{\nabla}g = \tilde{\nabla}\tilde{R} = \tilde{\nabla}h = \tilde{\nabla}S = 0, S(\tilde{R}) = \tilde{R}, S(h) = h, S(g) = g$ , where  $h = \nabla - \tilde{\nabla}$  and  $\tilde{R}$  is the curvature tensor field of  $\tilde{\nabla}$ .

Let  $M$  be such a  $k$ -s. l. R. s. that  $S_x = (s_x)_{*x}$  has only complex eigenvalues  $a_1 \pm b_1 i, \dots, a_r \pm b_r i$ . We define distributions  $D_i, i = 1, \dots, r$  by

$$D_i = \ker(S^2 - 2a_i S + I).$$

Every  $X \in \chi(M)$  has the unique decomposition  $X = X_1 + \dots + X_r$ , where  $X_i \in D_i, i = 1, \dots, r$ . An almost complex structure  $J$  on  $M$  is defined by

$$JX = \sum_{i=1}^r \frac{1}{b_i} (S - a_i I) X_i. \tag{6}$$

It is proved in [2] that  $(J, g)$  is an almost Hermitian structure and the connections  $\tilde{\nabla}$  and  $\bar{\nabla}$  defined by (1) and (5) respectively coincide i.e.  $\tilde{\nabla} = \bar{\nabla}$  on  $M$ .

**Proposition 1.5.** *Let  $(J_1, J_2, J_3, g)$  be such an ahHs on a  $k$ -s. l. R. s.  $(M, g)$  that  $J_1 = J$ , where  $J$  is defined by (6). In this case  $J_2$  and  $J_3$  are not invariant with respect to the family  $\{s_x: x \in M\}$ .*

**Proof.** Otherwise we have  $(s_x)_{*x} \cdot J_2 X = J_2 \cdot (s_x)_{*x} X$  i.e.  $S \cdot J_2 X = J_2 \cdot S X$ . Using (6) we get

$$J_2 J_1 X = \sum_{i=1}^r \frac{1}{b_i} (J_2 S - a_i J_2 X_i) = \sum_{i=1}^r \frac{1}{b_i} (S - a_i I) J_2 X_i,$$

$$J_1 J_2 X = \sum_{i=1}^r \frac{1}{b_i} (S - a_i I) (J_2 X)_i,$$

$$(S^2 - 2a_i S + I) J_2 X_i = J_2 (S^2 - 2a_i S + I) X_i = 0,$$

therefore  $J_2 X_i = (J_2 X)_i$  and  $J_1 J_2 X = J_2 J_1 X$  for any  $X \in \chi(M)$ . But we have  $J_1 J_2 X = - J_2 J_1 X$ , hence  $J_3 X = J_1 J_2 X = 0$  i.e.  $J_3 = 0$ .

We obtained the contradiction because  $J_3$  is nonsingular. By similar arguments  $J_3$  can not be invariant with respect to the family  $\{s_x: x \in M\}$ .  $\square$

## 2 HyperHermitian structures on tangent bundles

$0^0$ . Let  $(M, g)$  be a Riemannian manifold and  $TM$  be its tangent bundle. For a metric connection  $\tilde{\nabla}$  ( $\tilde{\nabla}g = 0$ ) we consider the connection map  $\tilde{K}$  of  $\tilde{\nabla}$  [1], [4], defined by the formula

$$\tilde{\nabla}_X Z = \tilde{K} Z_* X, \tag{7}$$

where  $Z$  is considered as a map from  $M$  into  $TM$  and the right side means a vector field on  $M$  assigning to  $p \in M$  the vector  $\tilde{K} Z_* X_p \in M_p$ .

If  $U \in TM$ , we denote by  $H_U$  the kernel of  $\tilde{K}|_{TM_U}$  and this  $n$ -dimensional subspace of  $TM_U$  is called the horizontal subspace of  $TM_U$ .

Let  $\pi$  denote the natural projection of  $TM$  onto  $M$ , then  $\pi_*$  is a  $C^\infty$ -map of  $TTM$  onto  $TM$ . If  $U \in TM$ , we denote by  $V_U$  the kernel of  $\pi_*|_{TM_U}$  and this  $n$ -dimensional subspace of  $TM_U$  is called the vertical subspace of  $TM_U$  ( $\dim TM_U = 2\dim M = 2n$ ). The following maps are isomorphisms of corresponding vector spaces ( $p = \pi(U)$ )

$$\pi_*|_{TM_U} : H_U \rightarrow M_p, \quad \tilde{K}|_{TM_U} : V_U \rightarrow M_p$$

and we have

$$TM_U = H_U \oplus V_U.$$

If  $X \in \chi(M)$ , then there exists exactly one vector field on  $TM$  called the “horizontal lift” (resp. “vertical lift”) of  $X$  and denoted by  $\bar{X}^h$  ( $\bar{X}^v$ ), such that for all  $U \in TM$ :

$$\pi_* \bar{X}_U^h = X_{\pi(U)}, \quad \tilde{K} \bar{X}_U^h = 0_{\pi(U)}, \tag{8}$$

$$\pi_* \bar{X}_U^v = 0_{\pi(U)}, \quad \tilde{K} \bar{X}_U^v = X_{\pi(U)}. \tag{9}$$

Let  $\tilde{R}$  be the curvature tensor field of  $\tilde{\nabla}$ , then following [1] we write

$$[\bar{X}^v, \bar{Y}^v] = 0, \tag{10}$$

$$[\bar{X}^h, \bar{Y}^v] = \left( \tilde{\nabla}_X \bar{Y} \right)^v, \tag{11}$$

$$\pi_*([\bar{X}^h, \bar{Y}^h]_U) = [X, Y], \tag{12}$$

$$\tilde{K}([\bar{X}^h, \bar{Y}^h]_U) = \tilde{R}(X, Y)U. \tag{13}$$

For vector fields  $\bar{X} = \bar{X}^h \oplus \bar{X}^v$  and  $\bar{Y} = \bar{Y}^h \oplus \bar{Y}^v$  on  $TM$  the natural Riemannian metric  $\langle , \rangle$  is defined on  $TM$  by the formula

$$\langle \bar{X}, \bar{Y} \rangle = g(\pi_*\bar{X}, \pi_*\bar{Y}) + g(\tilde{K}\bar{X}, \tilde{K}\bar{Y}). \tag{14}$$

It is clear that the subspaces  $H_U$  and  $V_U$  are orthogonal with respect to  $\langle , \rangle$ .

It is easy to verify that  $\bar{X}_1^h, \bar{X}_2^h, \dots, \bar{X}_n^h, \bar{X}_1^v, \bar{X}_2^v, \dots, \bar{X}_n^v$  are orthonormal vector fields on  $TM$  if  $X_1, X_2, \dots, X_n$  are those on  $M$  i.e.  $g(X_i, X_j) = \delta_j^i$ .

1<sup>0</sup>. We define a tensor field  $J_1$  on  $TM$  by the equalities

$$J_1\bar{X}^h = \bar{X}^v, J_1\bar{X}^v = -\bar{X}^h, X \in \chi(M). \tag{15}$$

For  $X \in \chi(M)$  we get

$$J_1^2\bar{X} = J_1(J_1(\bar{X}^h \oplus \bar{X}^v)) = J_1(-\bar{X}^h \oplus \bar{X}^v) = -(\bar{X}^h \oplus \bar{X}^v) = -I\bar{X}$$

and

$$J_1^2 = -I.$$

For  $X, Y \in \chi(M)$  we obtain

$$\begin{aligned} \langle J_1\bar{X}, J_1\bar{Y} \rangle &= \langle -\bar{X}^h \oplus \bar{X}^v, -\bar{Y}^h \oplus \bar{Y}^v \rangle = \langle -\bar{X}^h, -\bar{Y}^h \rangle + \langle \bar{X}^v, \bar{Y}^v \rangle, \\ \langle \bar{X}, \bar{Y} \rangle &= \langle \bar{X}^h \oplus \bar{X}^v, \bar{Y}^h \oplus \bar{Y}^v \rangle = \langle \bar{X}^h, \bar{Y}^h \rangle + \langle \bar{X}^v, \bar{Y}^v \rangle \end{aligned}$$

and it follows that  $\langle J_1\bar{X}, J_1\bar{Y} \rangle = \langle \bar{X}, \bar{Y} \rangle$ ,  $(TM, J_1, \langle , \rangle)$  is an almost Hermitian manifold.

Further, we want to analyze the second fundamental tensor field  $h^1$  of the pair  $(J_1, \langle , \rangle)$ .

The Riemannian connection  $\hat{\nabla}$  of the metric  $\langle , \rangle$  on  $TM$  is defined by the formula (see [4])

$$\begin{aligned} \langle \hat{\nabla}_{\bar{X}}\bar{Y}, \bar{Z} \rangle &= \frac{1}{2}(\bar{X} \langle \bar{Y}, \bar{Z} \rangle + \bar{Y} \langle \bar{Z}, \bar{X} \rangle - \bar{Z} \langle \bar{X}, \bar{Y} \rangle + \langle \bar{Z}, [\bar{X}, \bar{Y}] \rangle + \\ &+ \langle \bar{Y}, [\bar{Z}, \bar{X}] \rangle + \langle \bar{X}, [\bar{Z}, \bar{Y}] \rangle), X, Y, Z \in \chi(M). \end{aligned} \tag{16}$$

Using (2), (3) for orthonormal vector fields  $\bar{X}, \bar{Y}, \bar{Z}$  on  $TM$  we obtain

$$\begin{aligned} h^1_{\bar{X}\bar{Y}\bar{Z}} &= \langle h^1_{\bar{X}}\bar{Y}, \bar{Z} \rangle = \frac{1}{2} \langle \hat{\nabla}_{\bar{X}}\bar{Y} + J_1\hat{\nabla}_{\bar{X}}J_1\bar{Y}, \bar{Z} \rangle = \\ &= \frac{1}{2}(\langle \hat{\nabla}_{\bar{X}}\bar{Y}, \bar{Z} \rangle - \langle \hat{\nabla}_{\bar{X}}J_1\bar{Y}, J_1\bar{Z} \rangle) = \\ &= \frac{1}{4}(\langle [\bar{X}, \bar{Y}], \bar{Z} \rangle + \langle [\bar{Z}, \bar{X}], \bar{Y} \rangle + \langle [\bar{Z}, \bar{Y}], \bar{X} \rangle - \\ &- \langle [\bar{X}, J_1\bar{Y}], J_1\bar{Z} \rangle - \langle [J_1\bar{Z}, \bar{X}], J_1\bar{Y} \rangle - \langle [J_1\bar{Z}, J_1\bar{Y}], \bar{X} \rangle). \end{aligned} \tag{17}$$

Using (10)–(13) and (17) we consider the following cases for the tensor field  $h^1$  assuming all the vector fields to be orthonormal.

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For  $X, Y \in \chi(M)$  we get

$$\begin{aligned} \langle J_2\bar{X}, J_2\bar{Y} \rangle &= \langle (\overline{JX})^h \oplus -(\overline{JX})^v, (\overline{JY})^h \oplus -(\overline{JY})^v \rangle = \langle (\overline{JX})^h, (\overline{JY})^h \rangle + \\ &\quad + \langle (\overline{JX})^v, (\overline{JY})^v \rangle = g(JX, JY) + g(JX, JY) = \\ &= g(X, Y) + g(X, Y) = \langle \bar{X}^h, \bar{Y}^h \rangle + \langle \bar{X}^v, \bar{Y}^v \rangle \\ &= \langle \bar{X}^h \oplus \bar{X}^v, \bar{Y}^h \oplus \bar{Y}^v \rangle = \langle \bar{X}, \bar{Y} \rangle. \end{aligned}$$

Further, we obtain

$$\begin{aligned} J_1(J_2\bar{X}) &= J_1((\overline{JX})^h \oplus -(\overline{JX})^v) = (\overline{JX})^h \oplus (\overline{JX})^v, \\ J_2(J_1\bar{X}) &= J_2(-\bar{X}^h \oplus \bar{X}^v) = -(\overline{JX})^h \oplus -(\overline{JX})^v. \end{aligned}$$

Thus, we get  $J_1J_2 = -J_2J_1 = J_3$  and ahHs  $(J_1, J_2, J_3, \langle, \rangle)$  on  $TM$  has been constructed.

Using (2), (3) for orthonormal vector fields  $\bar{X}, \bar{Y}, \bar{Z}$  on  $TM$  we obtain

$$\begin{aligned} h^2_{\bar{X}\bar{Y}\bar{Z}} &= \langle h^2_{\bar{X}}\bar{Y}, \bar{Z} \rangle = \frac{1}{2} \langle \hat{\nabla}_{\bar{X}}\bar{Y} + J_2\hat{\nabla}_{\bar{X}}J_2\bar{Y}, \bar{Z} \rangle = \\ &= \frac{1}{2} (\langle \hat{\nabla}_{\bar{X}}\bar{Y}, \bar{Z} \rangle - \langle \hat{\nabla}_{\bar{X}}J_2\bar{Y}, J_2\bar{Z} \rangle) = \\ &= \frac{1}{4} (\langle [\bar{X}, \bar{Y}], \bar{Z} \rangle + \langle [\bar{Z}, \bar{X}], \bar{Y} \rangle + \langle [\bar{Z}, \bar{Y}], \bar{X} \rangle - \\ &\quad - \langle [\bar{X}, J_2\bar{Y}], J_2\bar{Z} \rangle - \langle [J_2\bar{Z}, \bar{X}], J_2\bar{Y} \rangle - \langle [J_2\bar{Z}, J_2\bar{Y}], \bar{X} \rangle). \end{aligned} \tag{19}$$

Using (10)–(13) and (19) we consider the following cases for the tensor field  $h^2$  assuming all the vector fields to be orthonormal.

$$\begin{aligned} h^2_{\bar{X}^h\bar{Y}^h\bar{Z}^h} &= \frac{1}{4} (\langle [\bar{X}^h, \bar{Y}^h], \bar{Z}^h \rangle + \langle [\bar{Z}^h, \bar{X}^h], \bar{Y}^h \rangle + \langle [\bar{Z}^h, \bar{Y}^h], \bar{X}^h \rangle - \\ &\quad - \langle [\bar{X}^h, J_2\bar{Y}^h], J_2\bar{Z}^h \rangle - \langle [J_2\bar{Z}^h, \bar{X}^h], J_2\bar{Y}^h \rangle - \\ &\quad - \langle [J_2\bar{Z}^h, J_2\bar{Y}^h], \bar{X}^h \rangle) = \frac{1}{4} (g([X, Y], Z) + g([Z, X], Y) + \\ &\quad + g([Z, Y], X) - g([X, JY], JZ) - g([JZ, X], JY) - \\ &\quad - g([JZ, JY], X)) = \frac{1}{2} (g(\nabla_X Y, Z) - g(\nabla_X JY, JZ)) = h_{XYZ}. \end{aligned} \tag{1.2^0}$$

$$\begin{aligned} h^2_{\bar{X}^h\bar{Y}^h\bar{Z}^v} &= \frac{1}{4} (\langle [\bar{X}^h, \bar{Y}^h], \bar{Z}^v \rangle + \langle [\bar{Z}^v, \bar{X}^h], \bar{Y}^h \rangle + \langle [\bar{Z}^v, \bar{Y}^h], \bar{X}^h \rangle - \\ &\quad - \langle [\bar{X}^h, J_2\bar{Y}^h], J_2\bar{Z}^v \rangle - \langle [J_2\bar{Z}^v, \bar{X}^h], J_2\bar{Y}^h \rangle - \\ &\quad - \langle [J_2\bar{Z}^v, J_2\bar{Y}^h], \bar{X}^h \rangle) = \frac{1}{4} (g(\tilde{R}(X, Y)U, Z) + g(\tilde{R}(X, JY)U, JZ)) = \\ &= -\frac{1}{4} (g(\tilde{R}(X, Y)Z, U) + g(\tilde{R}(X, JY)JZ, U)). \end{aligned} \tag{2.2^0}$$

By similar arguments we obtain

$$h^2_{\bar{X}^h\bar{Y}^v\bar{Z}^h} = \frac{1}{4} (g(\tilde{R}(X, Z)Y, U) + g(\tilde{R}(X, JZ)JY, U)) \tag{3.2^0}$$



$$h_{\tilde{X}^v \tilde{Y}^h \tilde{Z}^h}^2 = -\frac{1}{4}(g(\tilde{R}(Z, Y)X, U) - g(\tilde{R}(JZ, JY)X, U)). \tag{4.2^0}$$

$$h_{\tilde{X}^v \tilde{Y}^v \tilde{Z}^v}^2 = 0. \tag{5.2^0}$$

$$h_{\tilde{X}^v \tilde{Y}^v \tilde{Z}^h}^2 = 0. \tag{6.2^0}$$

$$h_{\tilde{X}^v \tilde{Y}^h \tilde{Z}^v}^2 = 0. \tag{7.2^0}$$

$$h_{\tilde{X}^h \tilde{Y}^v \tilde{Z}^v}^2 = \frac{1}{2}(g(\tilde{\nabla}_X Y, Z) - g(\tilde{\nabla}_X JY, JZ)). \tag{8.2^0}$$

It is clear that the construction of the ahHs on  $TM$  strongly depends on the connection  $\tilde{\nabla}$  and we can obtain in this way an infinite dimensional set of ahHs.

**Theorem 2.1.** *Let  $(M, g, J)$  be an almost Hermitian manifold. Then there exists an infinite family of ahHs on  $TM$  (in particular, such structures can be constructed by the method above).*

**Corollary 2.2.** *Let  $(M, g)$  be a Riemannian manifold. Then there exists an infinite set of ahHs on  $TTM$ .*

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